

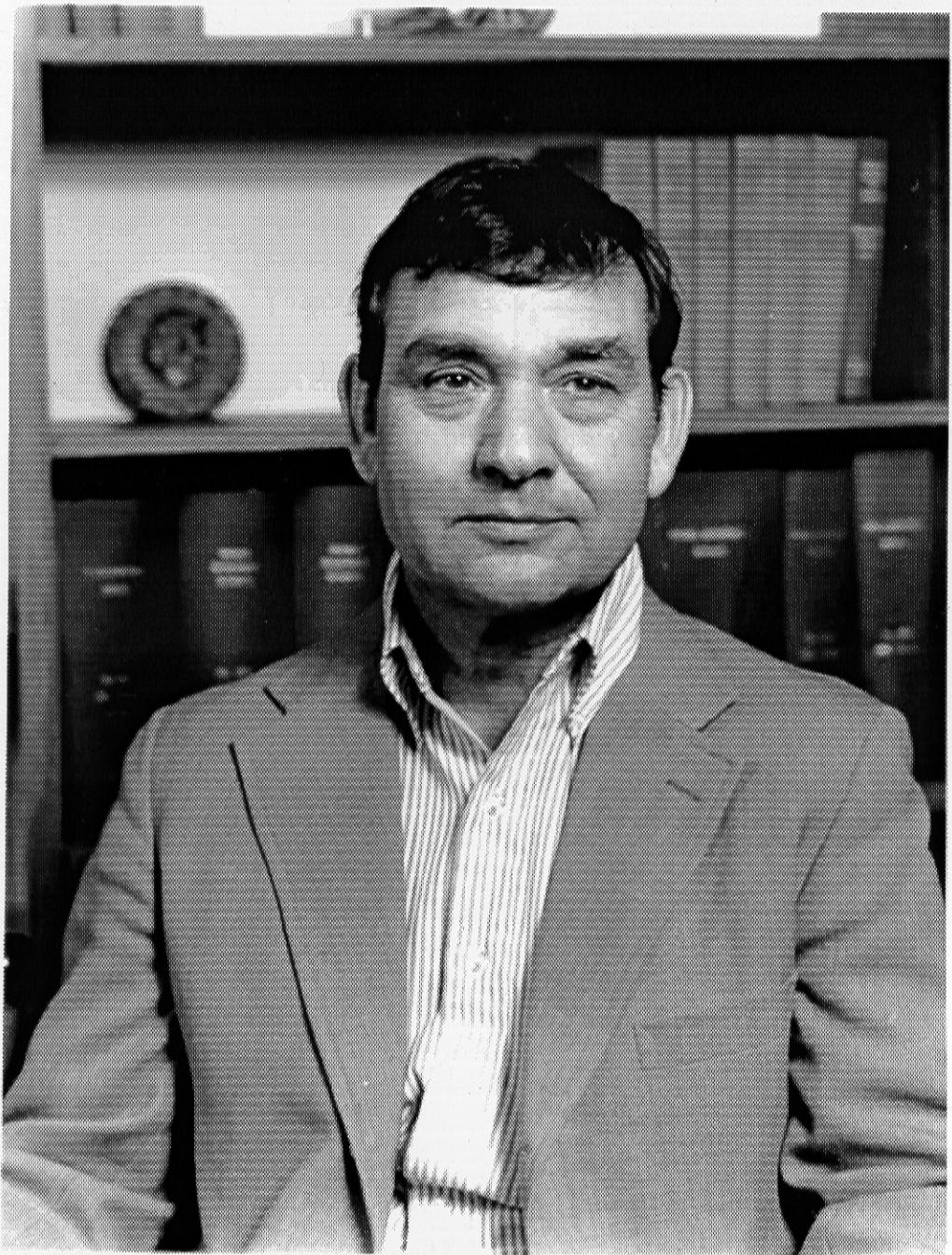
THE STATE OF THE STANDARD MODEL

CHRIS QUIGG

Fermilab

$\mu\mu^{99}$

San Francisco · 15 December 1999



Sam Treiman

ARNPS 46, 1 (1996)

OUR PICTURE OF MATTER

(November Revolution + 25 Years)

11.11.74

Pointlike ($r < 10^{-18}$ m) quarks

$$\begin{pmatrix} u \\ d' \end{pmatrix}_L \quad \begin{pmatrix} c \\ s' \end{pmatrix}_L \quad \begin{pmatrix} t \\ b' \end{pmatrix}_L$$

and leptons

$$\begin{pmatrix} \nu_e \\ e \end{pmatrix}_L \quad \begin{pmatrix} \nu_\mu \\ \mu \end{pmatrix}_L \quad \begin{pmatrix} \nu_\tau \\ \tau \end{pmatrix}_L$$

with interactions specified by

$$SU(3)_c \otimes SU(2)_L \otimes U(1)_Y$$

gauge symmetries

Q.C.D IS PART OF THE STD. MODEL...

A REMARKABLY
SIMPLE,
SUCCESSFUL,
AND RICH THEORY
(cf. F. Wilczek, hep-ph/9907340)

Perturbative QCD

- Exists, thanks to ASYMPTOTIC FREEDOM
- Describes many phenomena in quantitative detail:



Q^2 -evolution of Structure Fcns.



Jet production in $\bar{p}p$

► Many decays, evt. shapes, ...

We can measure the running of α_s
(engineering value for unification)

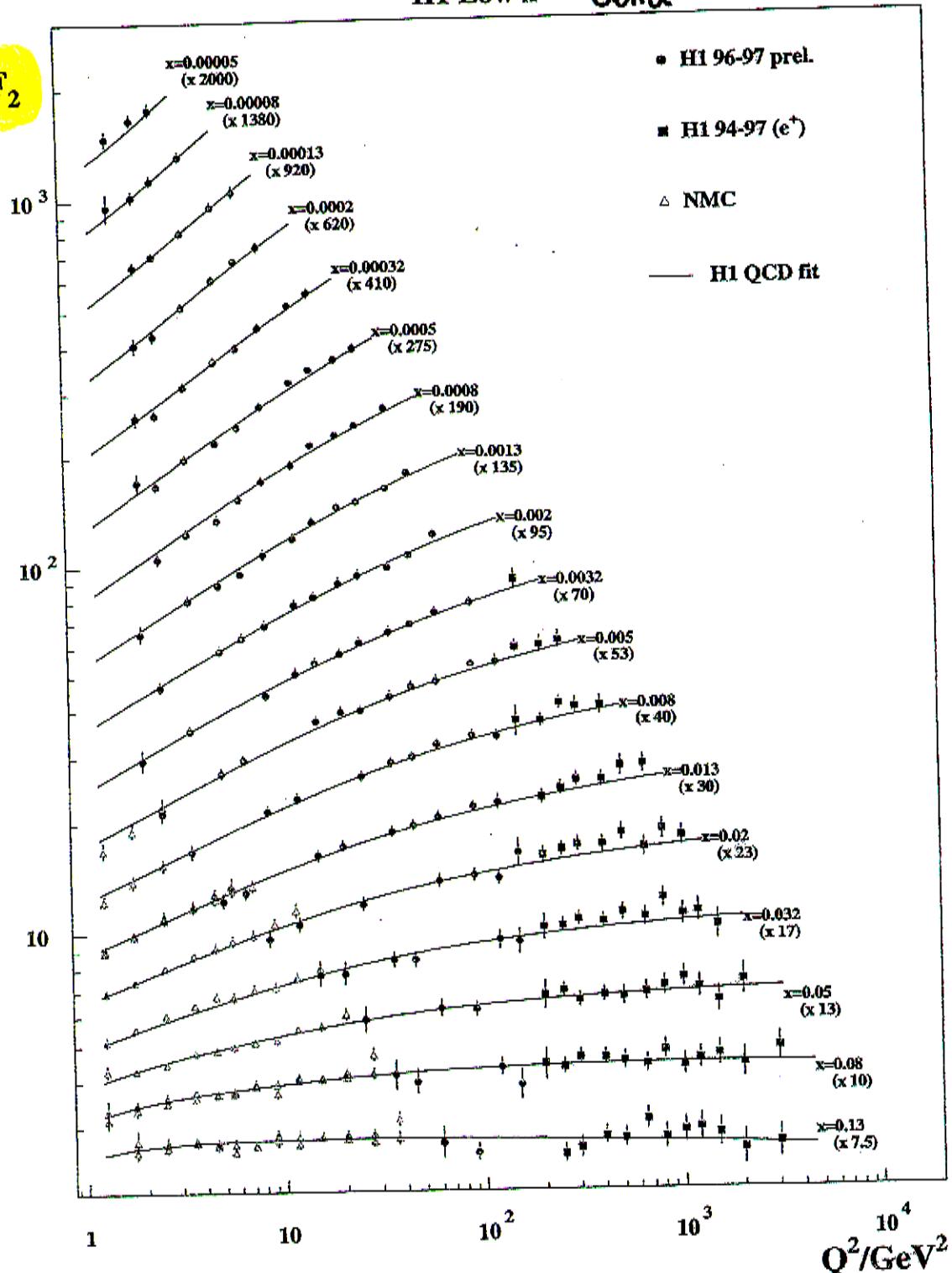
Nonperturbative (Lattice) QCD

- Predicts the hadron spectrum
- Gives our best information on quark masses, etc.

Large scaling violations described by NLO DGLAP

$$\frac{\partial F_2}{\partial \ln Q^2} \sim d_s \cdot x g$$

H1 Low x



where is BFKL , \ln^4/x terms ?

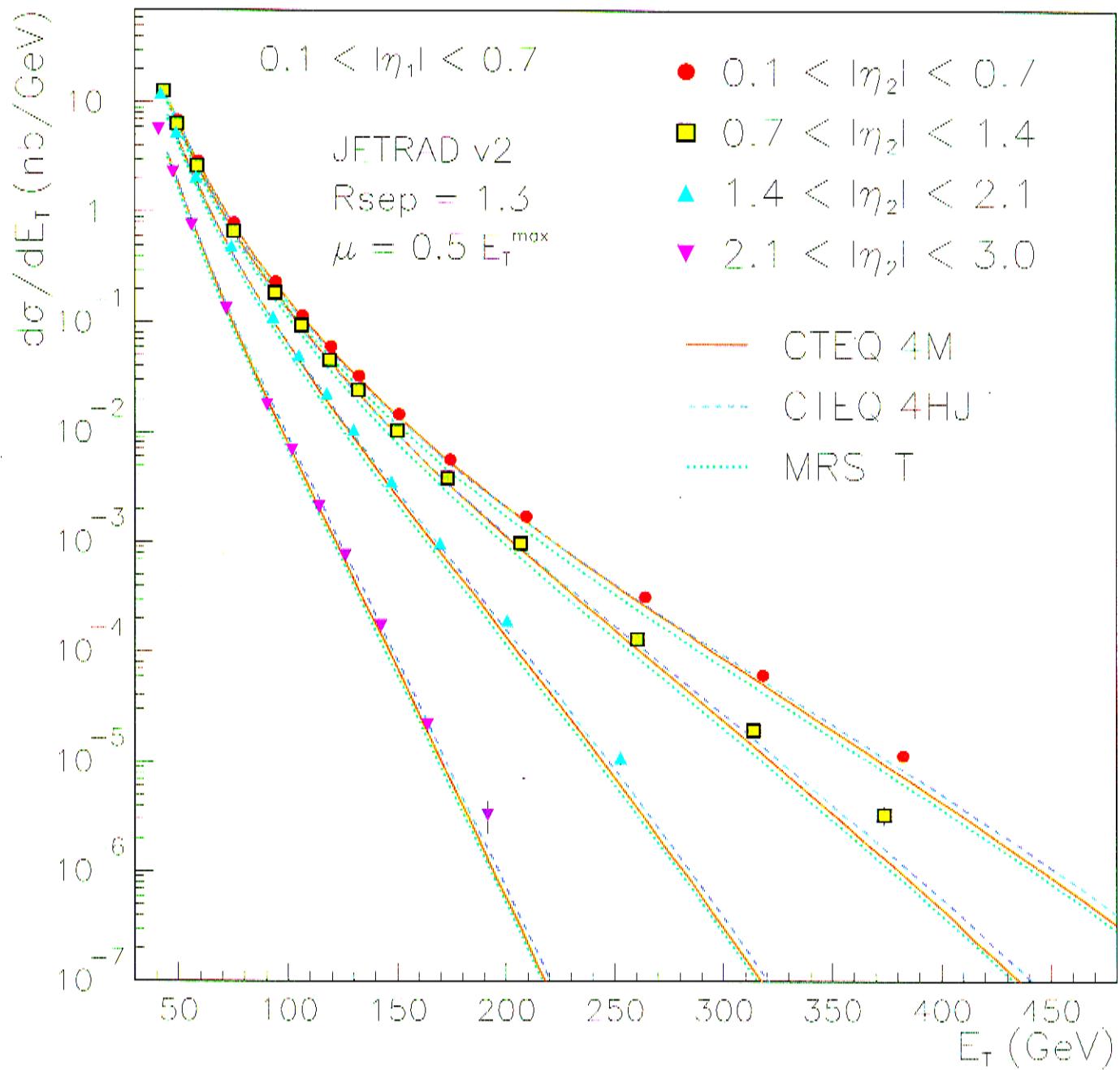
Balitskii, Fadin, Kuraev, Lipatov 1975,78

cf. Proc. DIS99

Grafaloni, Fadin, del Duca
Thorne ...

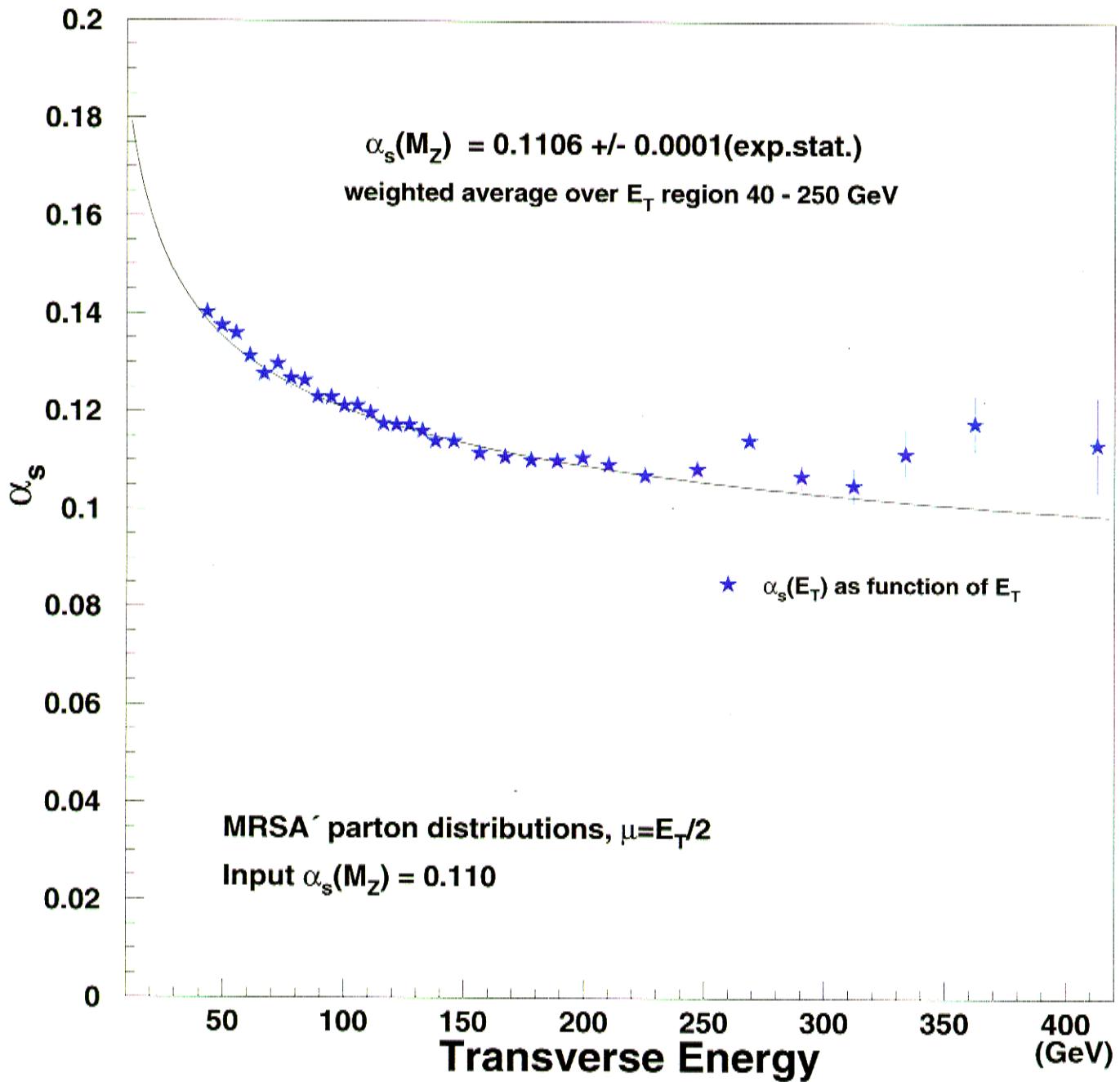
$\bar{p}p \rightarrow \text{jet}_1 + \text{jet}_2 + \text{anything}$

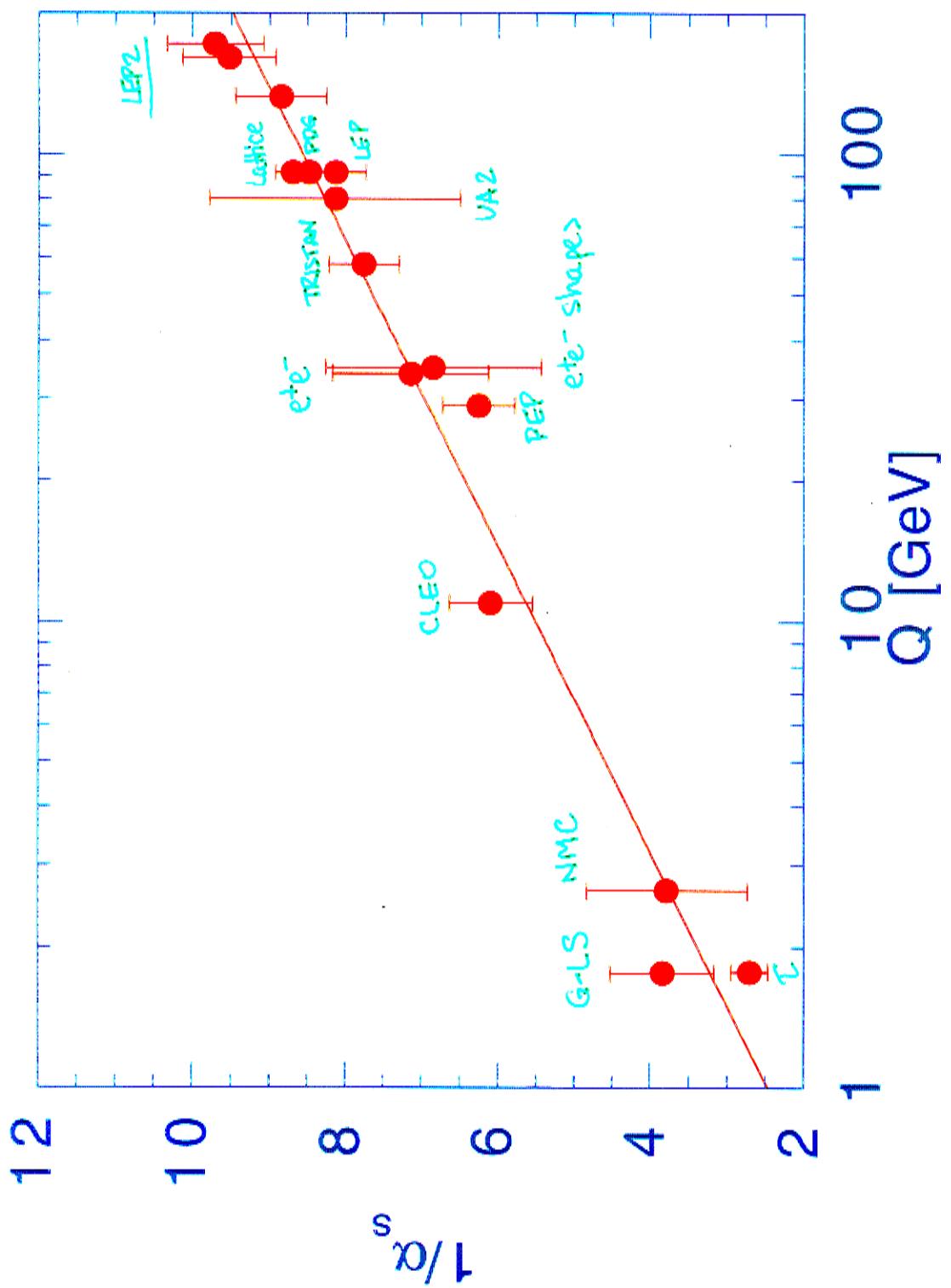
CDF Preliminary

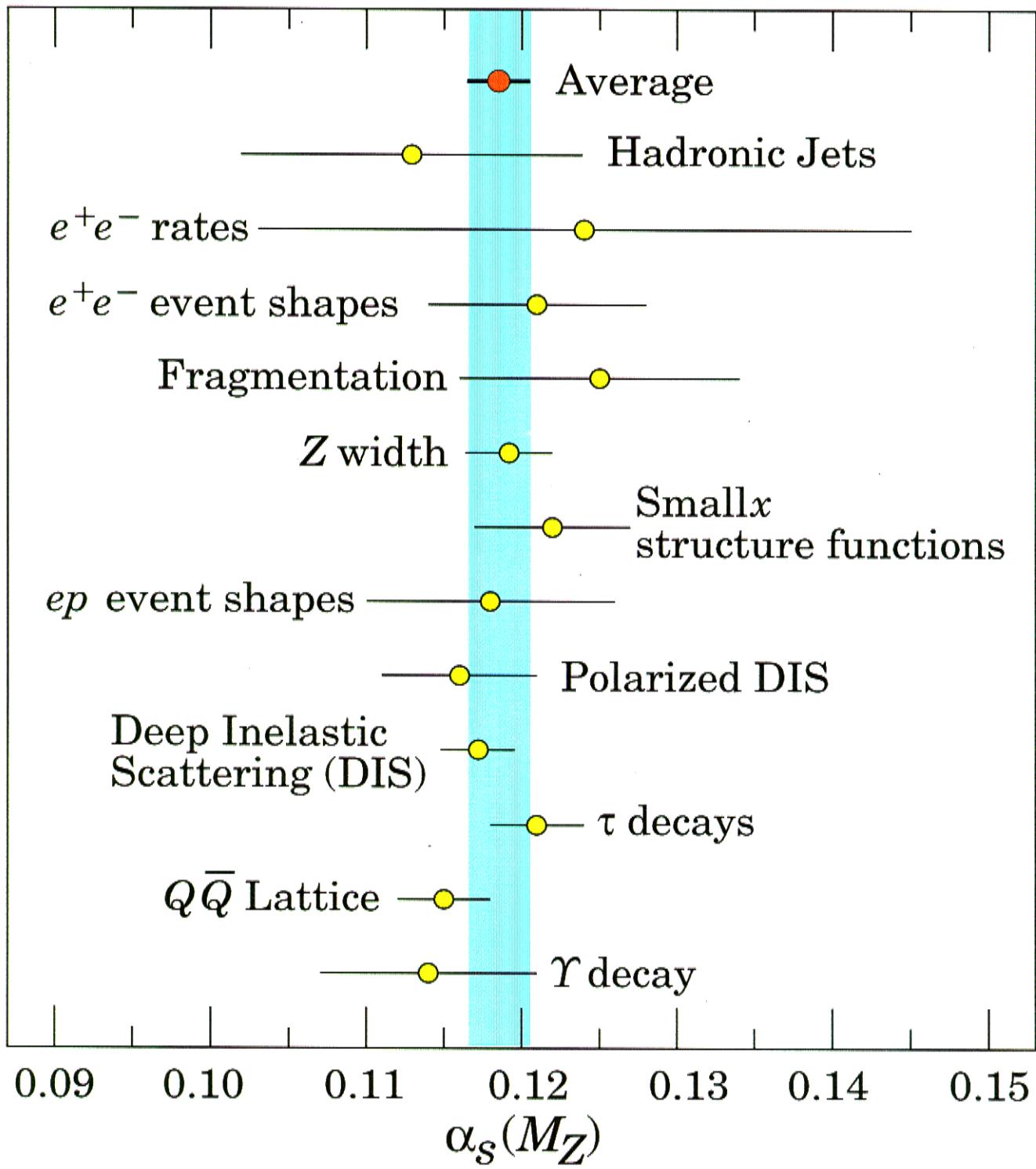


α_s determined from $\bar{p}p \rightarrow \text{jet}_1 + \text{jet}_2 + \text{anything}$

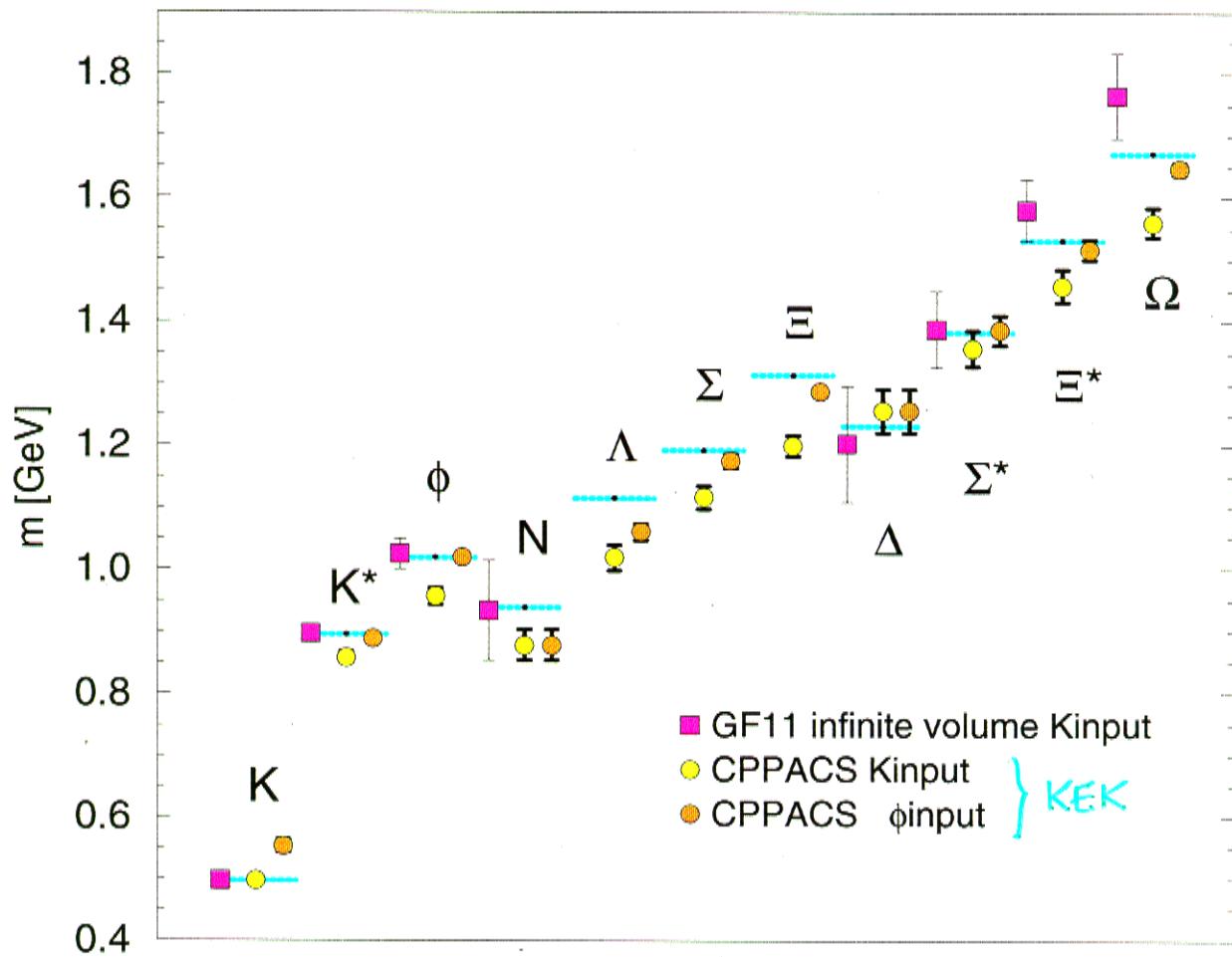
CDF Preliminary







PDG



No dynamical fermions (coming soon)

R. Burkhalter, hep-lat/9810043

THE $SU(2)_L \otimes U(1)_Y$ ELECTROWEAK THEORY

$$\begin{pmatrix} u \\ d' \end{pmatrix}_L \quad \begin{pmatrix} c \\ s' \end{pmatrix}_L \quad \begin{pmatrix} t \\ b' \end{pmatrix}_L$$

$$\begin{pmatrix} \nu_e \\ e \end{pmatrix}_L \quad \begin{pmatrix} \nu_\mu \\ \mu \end{pmatrix}_L \quad \begin{pmatrix} \nu_\tau \\ \tau \end{pmatrix}_L$$

+ IDEALIZATION THAT NEUTRINOS
ARE MASSLESS

- MANY SUCCESSES:

Neutral Currents

Charm

W^\pm, Z^0

..

- PRECISION MEASUREMENTS (10 years of LEPetal.)

Testing the Quantum Field Theory

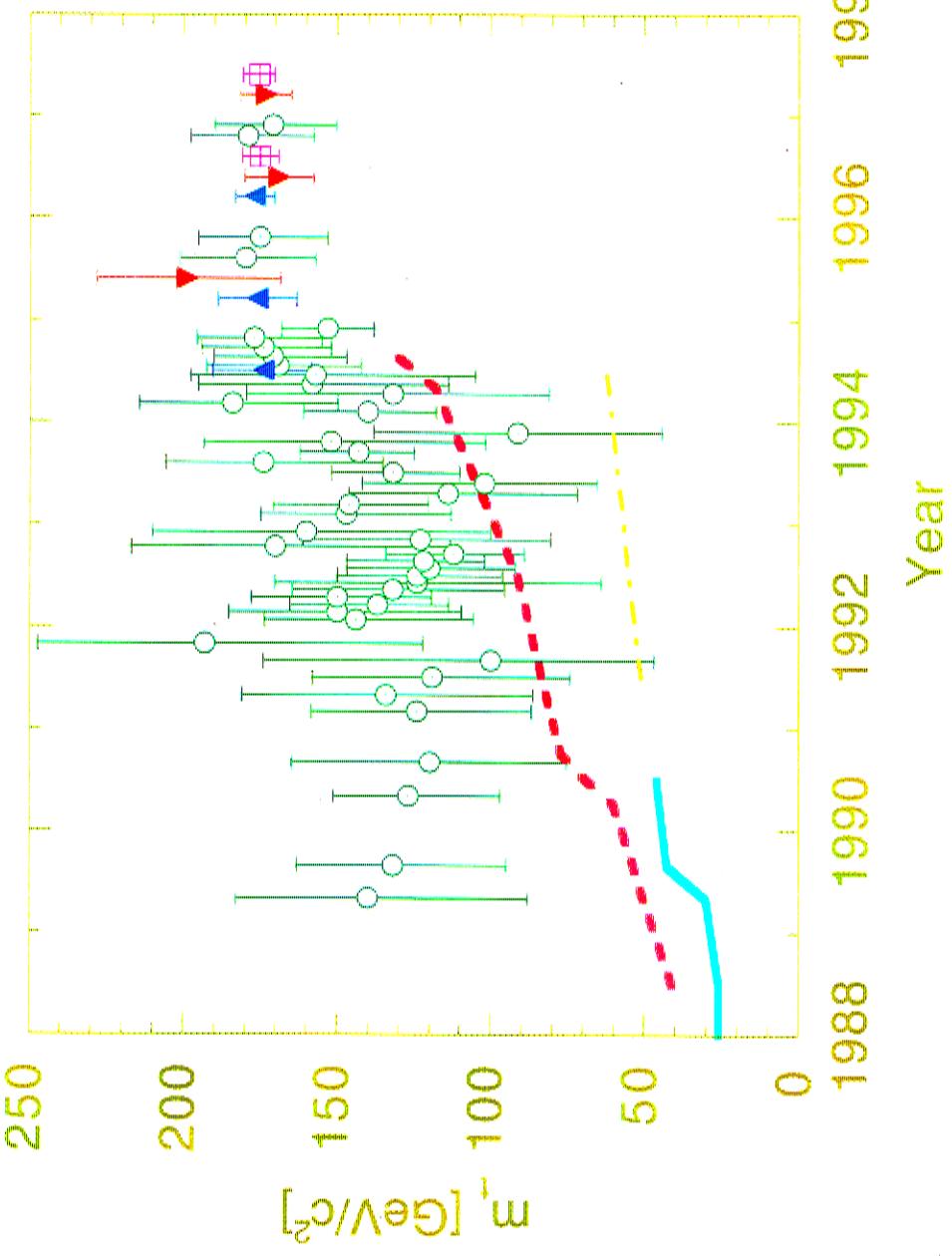
Looking for new physics

"in the sixth place of decimals"

Cf. A. Sirlin, hep-ph/9912227

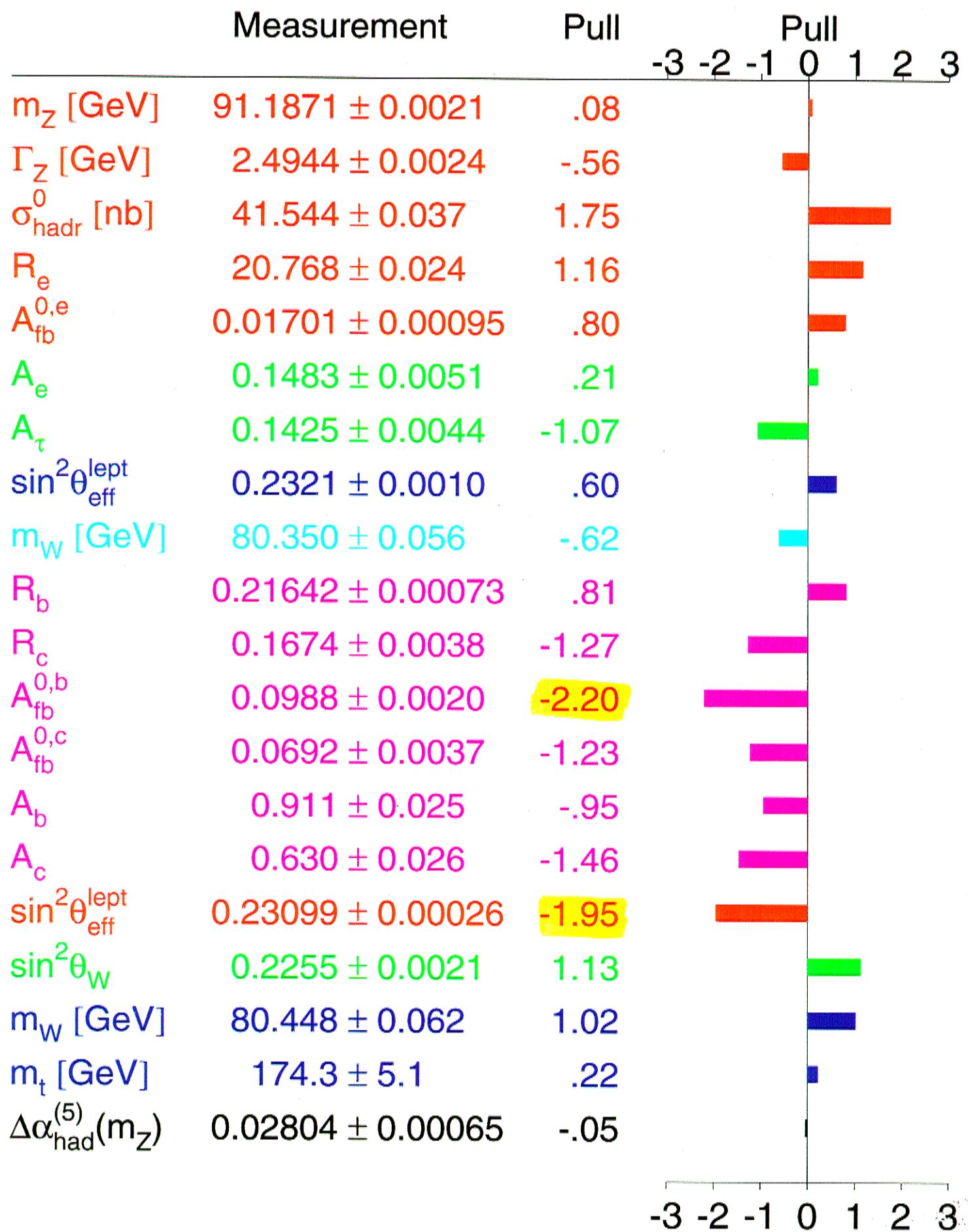
Global Fits to precision EW measurements

- precision improves with time
- calculations have improved with time



$$11.94: \text{LEP EW fit: } m_t = 178 \pm 11^{+18}_{-19} \text{ GeV}/c^2 \quad (174.3 \pm 5.1 \text{ GeV}/c^2)$$

Stanford 1999



Cho + Hagiwara, hep-ph/9912260

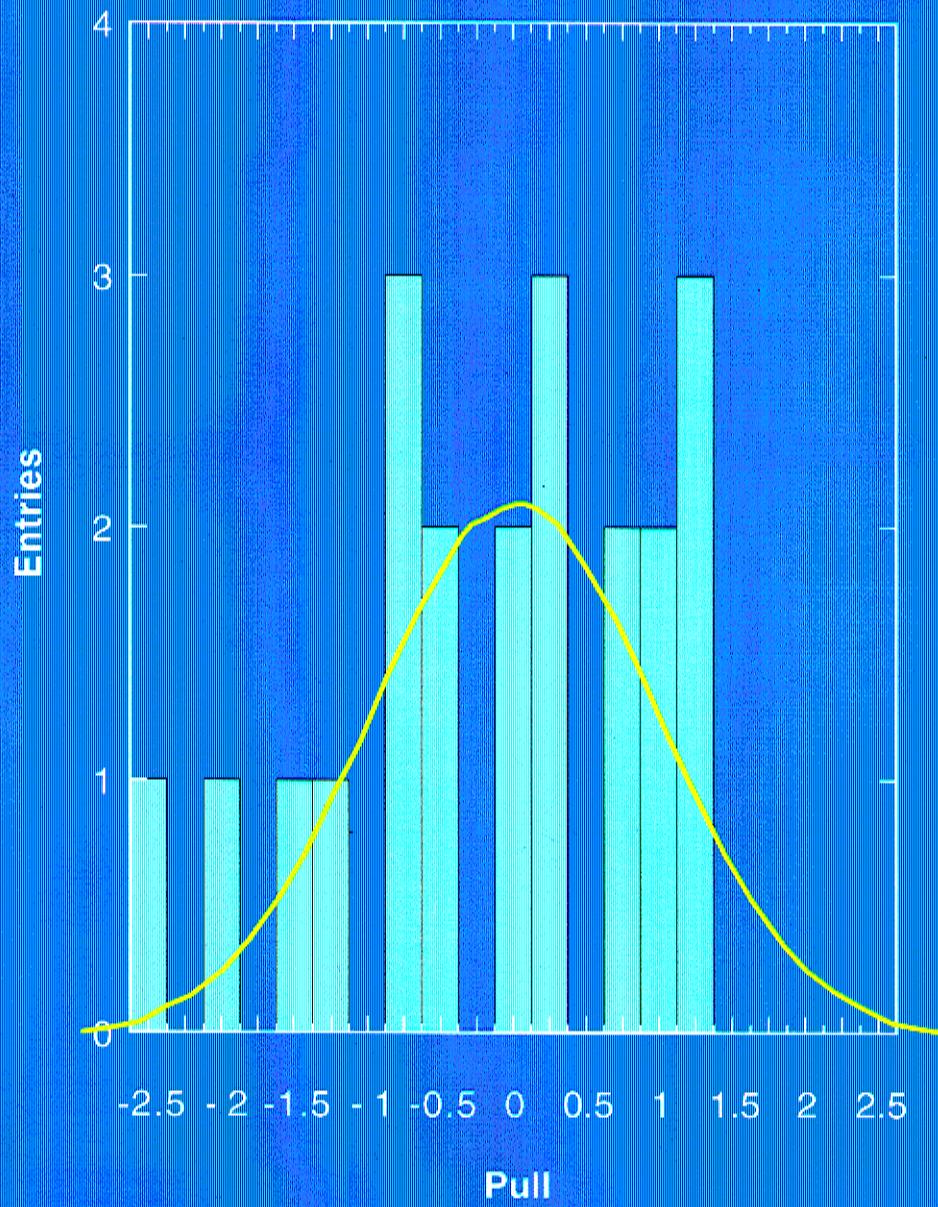
	data	SM*	pull*
LEP 1 [11]			
line-shape & FB asym.:			
m_Z (GeV)	91.1867 ± 0.0021	—	—
Γ_Z (GeV)	2.4939 ± 0.0024	2.4972	-1.4
σ_h^0 (nb)	41.491 ± 0.058	41.474	0.3
R_ℓ	20.765 ± 0.026	20.747	(0.7)
$A_{FB}^{0,\ell}$	0.01683 ± 0.00096	0.01651	(0.3)
for each lepton:			
$\left\{ \begin{array}{l} R_e \\ R_\mu \\ R_\tau \end{array} \right.$	20.783 ± 0.052 20.789 ± 0.034 20.764 ± 0.045	20.747 20.747 20.795	0.7 1.3 -0.7
$\left\{ \begin{array}{l} A_{FB}^{0,e} \\ A_{FB}^{0,\mu} \\ A_{FB}^{0,\tau} \end{array} \right.$	0.0153 ± 0.0025 0.0164 ± 0.0013 0.0183 ± 0.0017	0.0165 0.0165 0.0165	-0.5 -0.1 1.1
τ polarization:			
A_τ	0.1431 ± 0.0045	0.1484	-1.2
A_c	0.1479 ± 0.0051	0.1484	-0.1
b and c quark results:			
R_b	0.21656 ± 0.00074	0.21566	1.2
R_c	0.1735 ± 0.0044	0.1721	0.3
$A_{FB}^{0,b}$	0.0990 ± 0.0021	0.1040	-2.4
$A_{FB}^{0,c}$	0.0709 ± 0.0044	0.0744	-0.8
jet charge asymmetry:			
$\sin^2 \theta_{\text{eff}}^{\text{lept}}$	0.2321 ± 0.0010	0.2314	0.7
SLC [11]			
A_{LR}^0	0.1510 ± 0.0025	0.1484	1.0
A_b	0.867 ± 0.035	0.935	-1.9
A_c	0.647 ± 0.040	0.668	-0.5
Tevatron + LEP 2 [32]			
m_W (GeV)	80.410 ± 0.044	80.402	0.18
χ^2_{tot} (19 data points)			19.8
Parameters	Constraints		
m_t (GeV) [29]	173.8 ± 5.2	175.0	—
$\alpha_s(m_Z)$ [29]	0.119 ± 0.002	0.118	—
$1/\alpha(m_Z^2)$ [30]	128.90 ± 0.09	128.90	—
[31]	128.94 ± 0.04	—	—

Table 1: Electroweak measurements at LEP, SLC and Tevatron. The average W -boson mass is found in ref. [32]. The reference SM predictions and the corresponding ‘pull’ factors are given for $m_t = 175$ GeV, $m_{H_{\text{SM}}} = 100$ GeV, $\alpha_s(m_Z) = 0.118$ and $1/\alpha(m_Z^2) = 128.90$. Correlation matrix elements of the Z line-shape parameters and those for the heavy-quark parameters are found in ref. [11]. The data R_ℓ and $A_{FB}^{0,\ell}$ are obtained by assuming the $e\text{-}\mu\text{-}\tau$ universality and are not used in our χ^2 analysis.

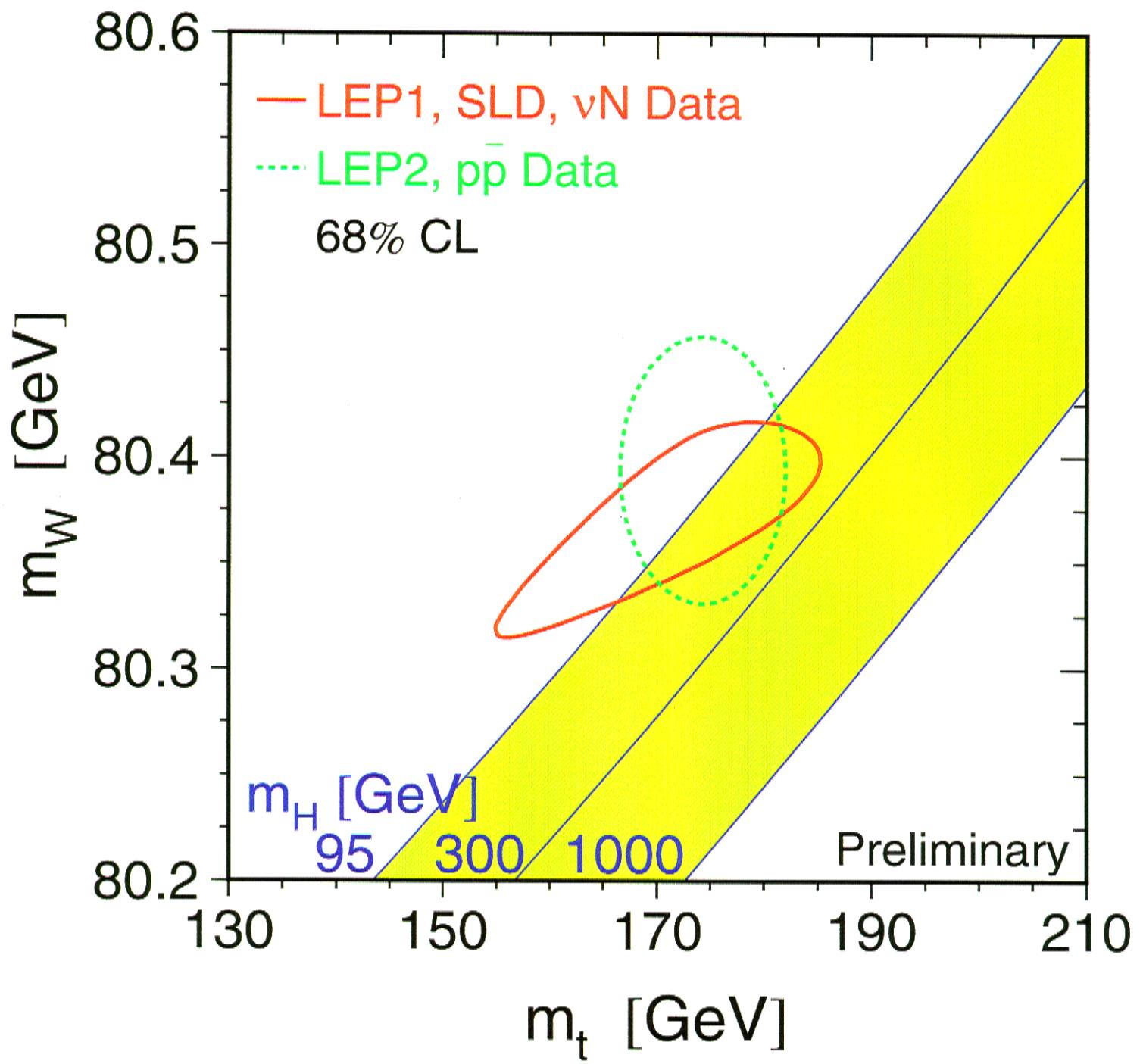
(cf. W. J. Marciano, hep-ph/9902332)

Anomalous
 $Z \rightarrow b_R \bar{b}_R$?
 ↗ FCNC in
 $b \rightarrow s, d; s \rightarrow d$

Standard Model Fit



Cho + Hagiwara



(Bill Marciano for details)

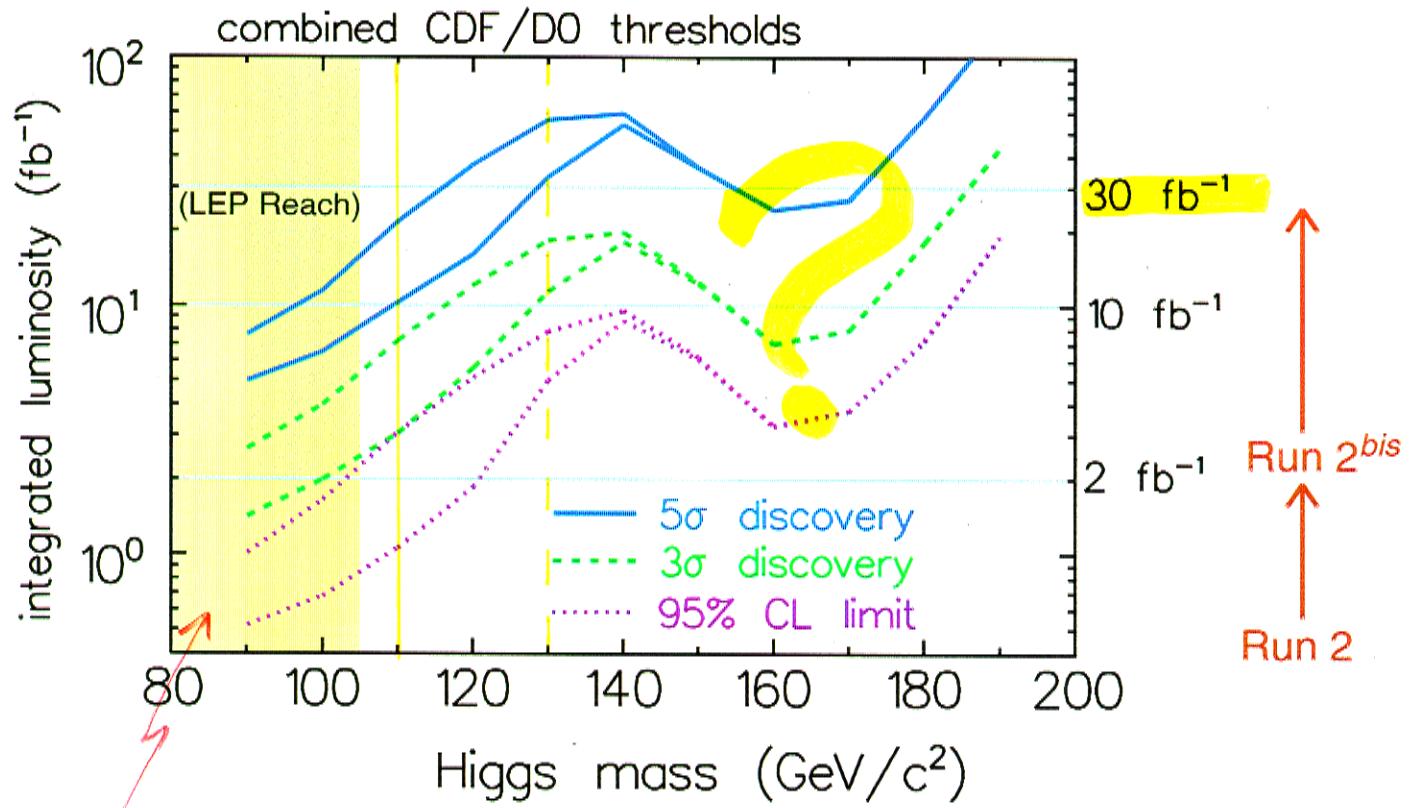
SEARCH FOR THE HIGGS BOSON AT THE TEVATRON

(J. Hobbs)

$$\bar{p}p \rightarrow H + (W^\pm, Z) + \dots$$

$\hookrightarrow b\bar{b}$

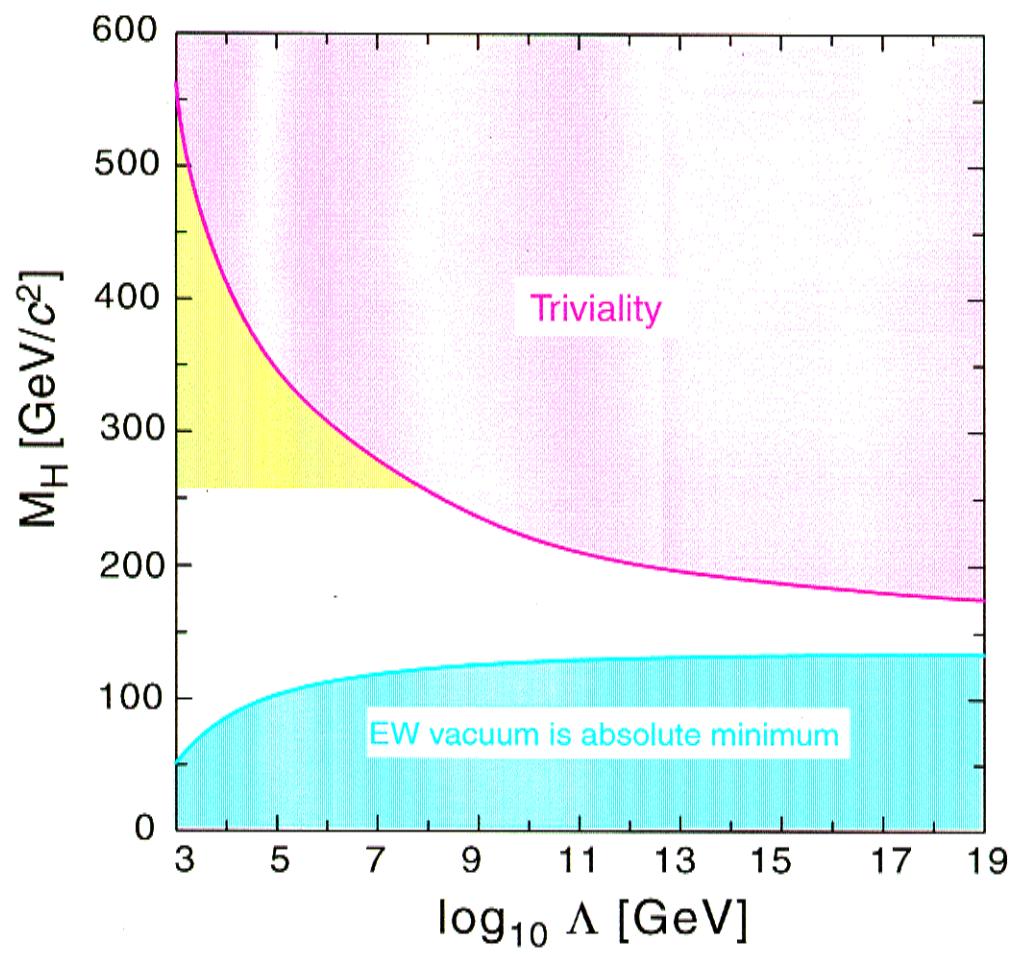
Key improvements: \mathcal{L} , 10% $b\bar{b}$ mass resolution,
measured in Run 2 $Z^0 \rightarrow b\bar{b}$



Sau-Lan Wu

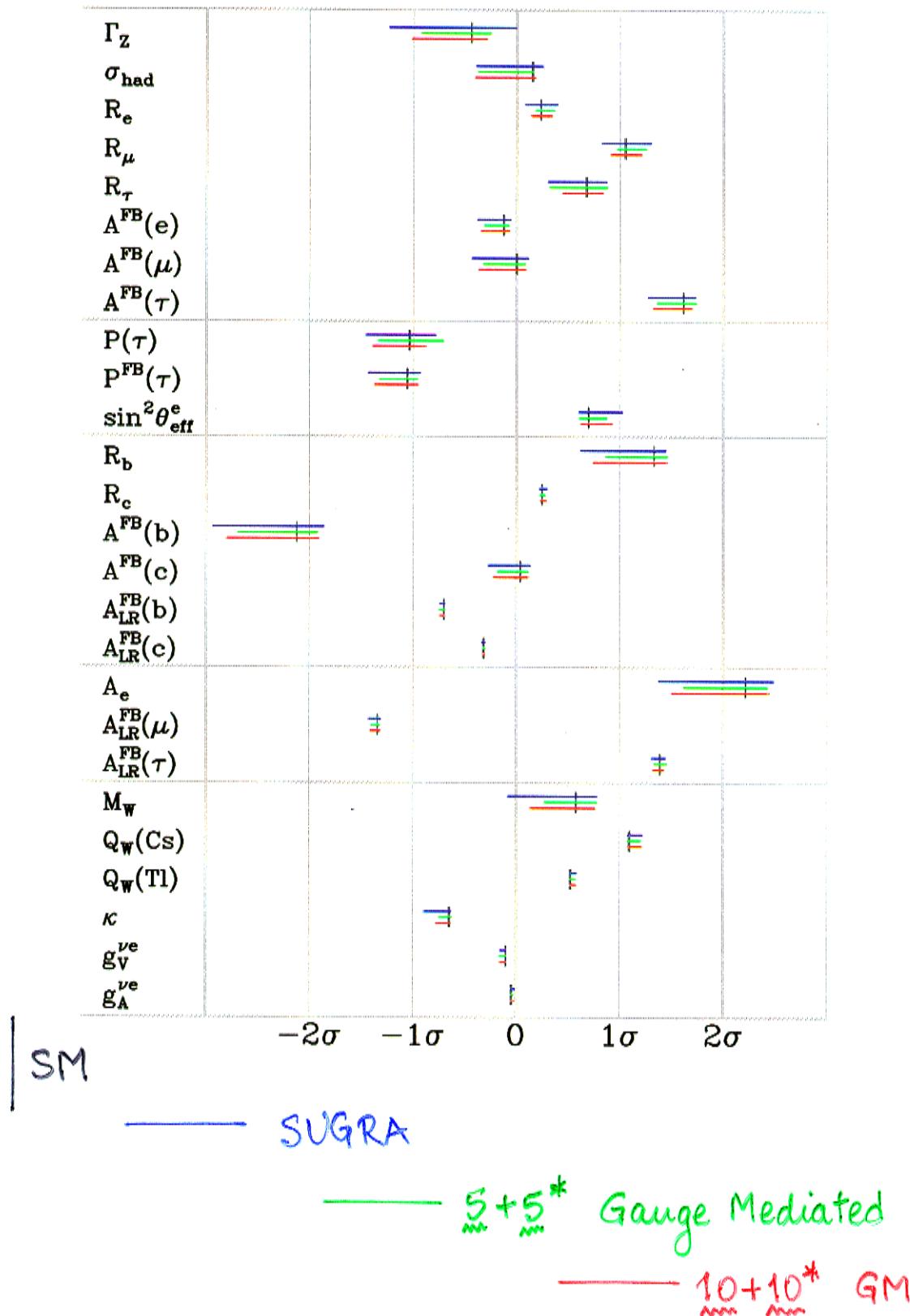
Goal of 30 fb^{-1} by 2007

Extend reach to $M_H \approx 180 \text{ GeV}/c^2$
using $H \rightarrow WW^*$?



MSSM Closely Resembles the Standard EW Theory

Eder + Pierce: SUSY vs. SM (hep-ph/9801238)



PRECISION EW CONSTRAINTS ON SUSY SPECTRUM?

(Many recent papers)

Cho + Hagiwara (hep-ph/9912260) best overall fit:

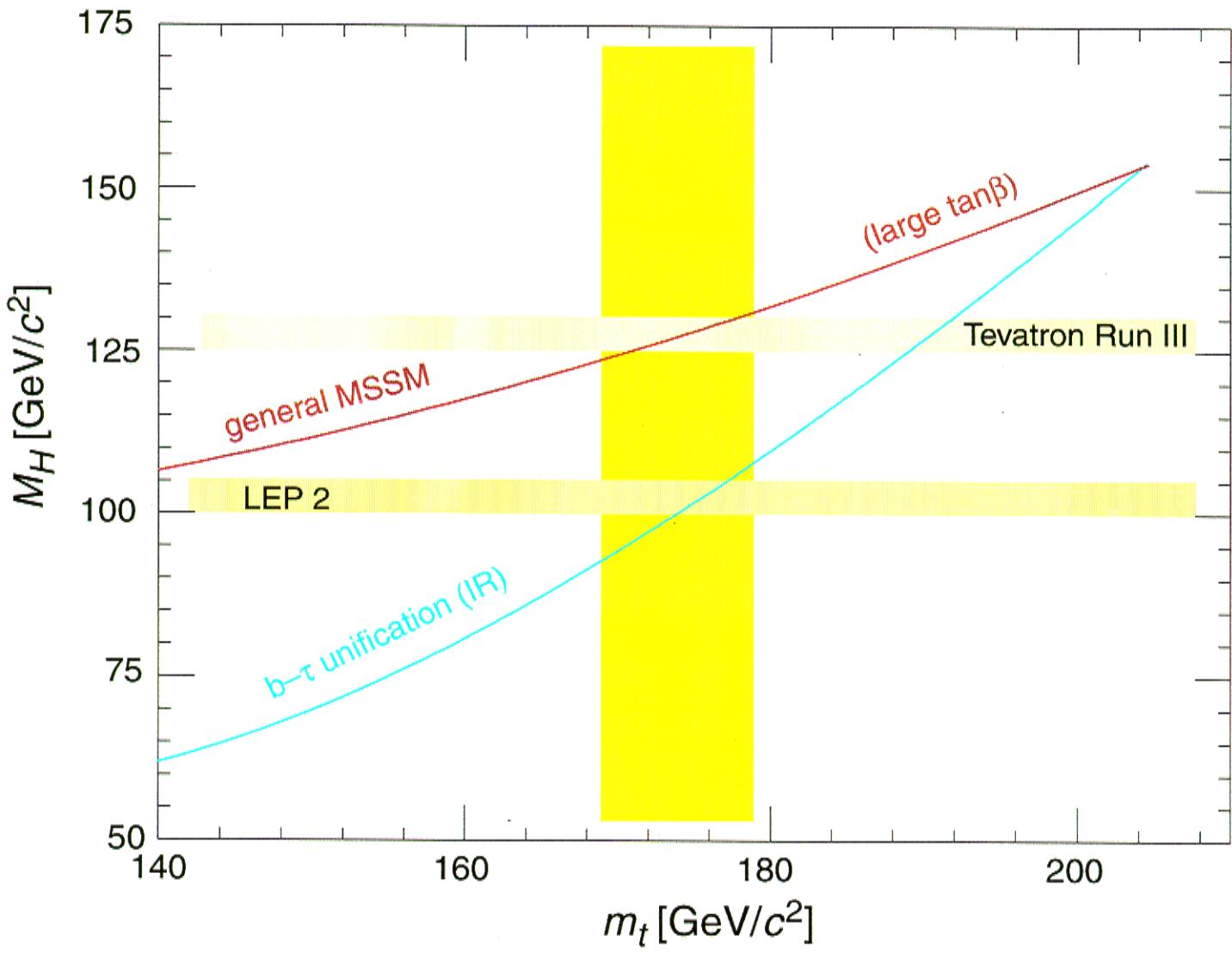
(WINO-LIKE) CHARGEDOS, $M \sim 100$ GeV

$$\tilde{\chi}_1^0 \quad \sim 50 \text{ GeV}$$

DOUBLET SQUARKS $\gtrsim 200$ GeV

SLEPTONS heavier

$$m_{\tilde{t}_1} \approx 134-169 \text{ GeV}$$



HOW MUCH CAN WE TRUST SM ANALYSIS TO BOUND M_H ?

- Hall and Kolda

Include higher-dimension ops.
suggested by extra space dim.

CONSTRAINTS EVAPORATE!

Weakness - not shown to be
the outcome of a real theory.

- Collins, Grant, Georgi hep-ph/9908330

Topcolor seesaw model

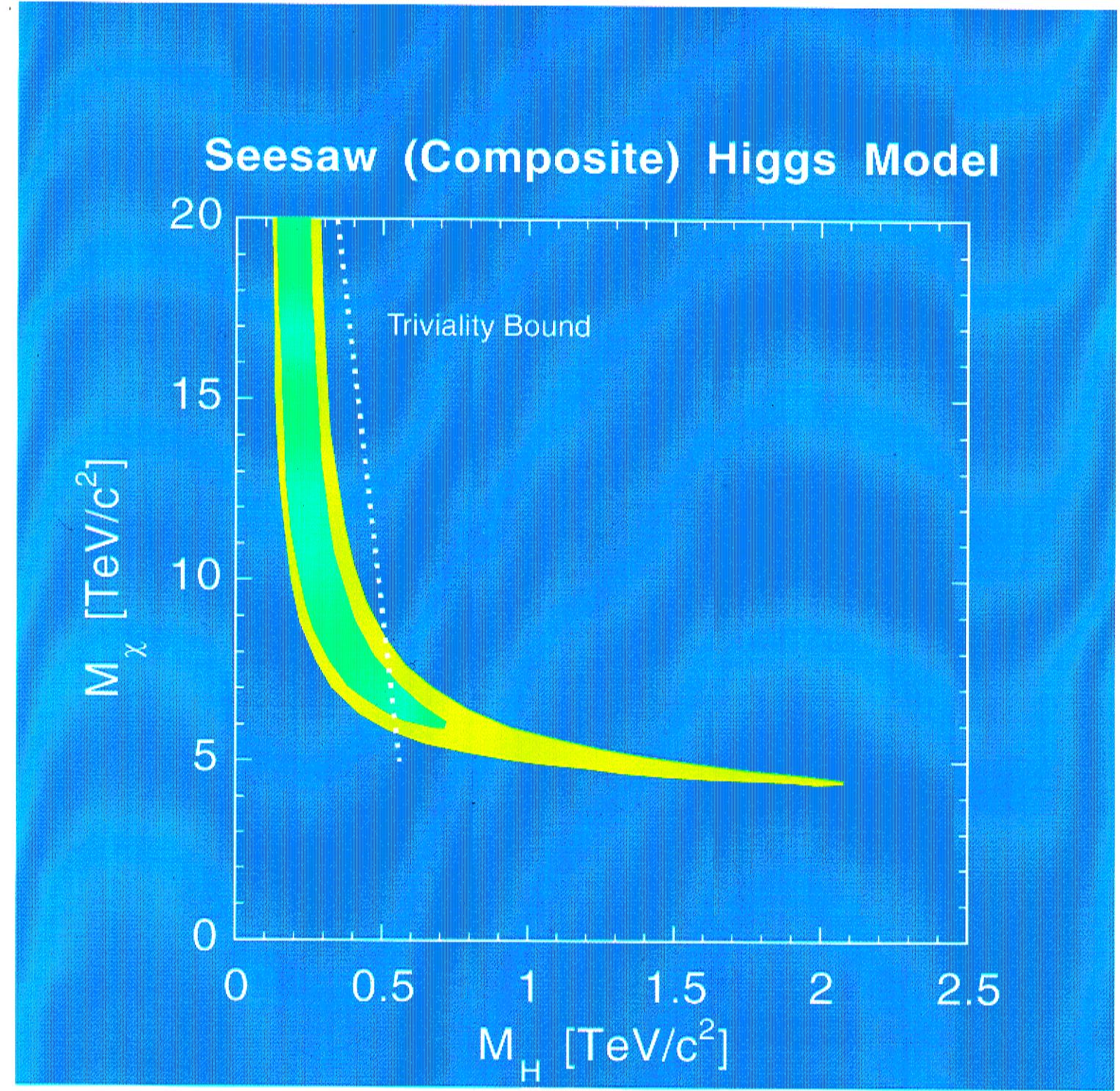
(1 additional heavy weak- $_m^1$
fermion X with $Y = 4/3$)

- $M_H \gtrsim 300 \text{ GeV}$ for $m_X \approx 5-7 \text{ TeV}$

(+ amusing model with 2 heavy fermions
+ 2 Higgs doublets \Rightarrow heavy h^0, H^0, H^\pm , but
light pseudoscalar)

See also: Chivukula + Evans, hep-ph/9907414

S. Chivukula's Circle Line talk (www-theory.fnal.gov)



LOOKING FOR TROUBLE

Bennett + Wieman PRL 82, 2484 (99)

measure weak charge Q_W in $6S-7S$ transition in Cs.

$$Q_W = -72.06 \pm 0.28 \text{ (exp.)} \pm 0.34 \text{ (th.)}$$

↗

improved

7-fold improvement

$$\approx Q_W^{(\text{SM})} + 2.5 \sigma$$

Enler + Langacker hep-ph/9910315

note that Q_W and Γ_Z^* can be improved by adding a Z' with mass ~ 800 GeV (E_6 Z_X)

* in their fit, $N_\nu = 2.985 \pm 0.008 = 3 - 2\sigma$

The Vacuum Energy Problem (Veltman)

Higgs potential

$$V(\varphi^+\varphi) = \mu(\varphi^+\varphi) + |\lambda|(\varphi^+\varphi)^2$$

At the minimum

$$\langle\varphi\rangle_0 = \begin{pmatrix} 0 \\ v/\sqrt{2} \end{pmatrix} \approx \begin{pmatrix} 0 \\ 176 \text{ GeV} \end{pmatrix},$$

$$V(\langle\varphi^+\varphi\rangle_0) = \frac{\mu^2 v^2}{4} = -\frac{|\lambda| v^4}{4} < 0$$

IDENTIFY $M_H^2 = -2\mu^2$.

Higgs potential contributes a field-independent constant energy density

$$\rho_H \equiv \frac{M_H^2 v^2}{8} \gtrsim \underbrace{10^8 \text{ GeV}^4}_{M_H \gtrsim 100 \text{ GeV}}$$

Observed Vacuum Energy Density

$$\rho_{\text{vac}} \lesssim 10^{-46} \text{ GeV}^4$$

MISMATCH BY **54 ORDERS OF MAG.**

How to zero ρ_H ?

THE PROBLEMS OF MASS, AND OF MASS SCALES

1) EWSB sets M_W, M_Z

$$M_W^2 = g^2 v^2 / 2 = \pi \alpha / G_F \sqrt{2} \sin^2 \theta_W$$

$$M_Z^2 = M_W^2 / \cos^2 \theta_W$$

$$\text{EW scale is } v = (G_F \sqrt{2})^{-1/2} = 246 \text{ GeV}$$

But it is not the only scale...

certain: $M_{\text{Planck}} = 1.22 \times 10^{19} \text{ GeV}$ (from G_N)

probable: $SU(3)_C \otimes SU(2)_L \otimes U(1)_Y$ unification scale

somewhere: flavor scale

→ Famous problem of Higgs scalar mass:

how to keep scales from mixing?

how to stabilize M_H ?

(unification doesn't help)

Problems of mass ...

2) Each fermion mass \Rightarrow a new, unknown, Yukawa coupling

$$L \equiv \begin{pmatrix} \nu \\ e \end{pmatrix}_L, \quad R \equiv e_R$$

$$\mathcal{L}_{\text{Yuk}}^{(e)} = -\xi_e [\bar{R}(\varphi^\dagger L) + (\bar{L}\varphi) R]$$

$$m_e = \xi_e v / \sqrt{2}$$

All fermion masses \sim physics beyond the S.M.

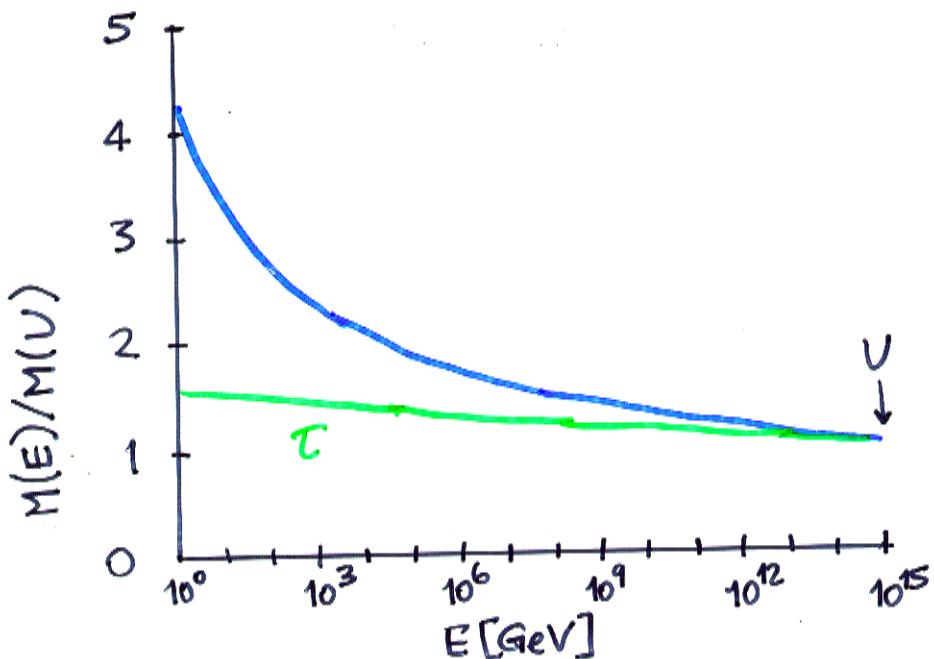
$$\xi_t \approx 1, \quad \xi_e \approx 3 \times 10^{-6}, \quad \xi_\nu \approx 10^{-10} ??$$

What accounts for the range (and values)
of the Yukawa couplings?

Best hope until now:

Unified theories suggest
pattern of fermion masses
simplifies on high scales

Famous triumph: m_b/m_τ



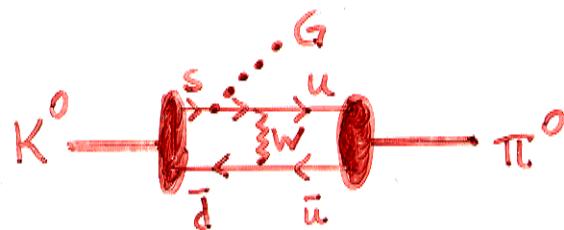
Perhaps all fermion masses arise on high scales
and show simple patterns there?

+ G_N small (M_{Pl} large) \Rightarrow

Gravitation irrelevant for particle physics

$$q \nearrow \begin{matrix} q \\ \dots G \end{matrix} \quad \sim \frac{E}{M_{Pl}}$$

Estimate $B(K \rightarrow \pi G) \sim \left(\frac{M_K}{M_{Pl}}\right)^2 \sim 10^{-38}$



+ We know from EW theory alone

that the 1-TeV scale is special

$$M_H^2 < \frac{8\pi\sqrt{2}}{3G_F} \simeq 1 \text{ TeV}^2 \quad (\text{Higgs boson or new physics})$$

So to eliminate the hierarchy problem,
consider EW physics beyond the standard model.

$$SU(3)_c \otimes SU(2)_L \otimes U(1)_Y$$

Composite Higgs
(Technicolor, ...)
Supersymmetry

CONVENTIONAL APPROACH TO NEW PHYSICS

Extend Standard Model to understand
why $(M_H, v) \ll M_{Pl}$ (A. Masiero)

NOVEL SPECULATION

Change Gravity to understand
why $M_{Pl} \gg v$

Expt.: Gravity follows Newtonian force law
down to 1 mm. (FIG.)

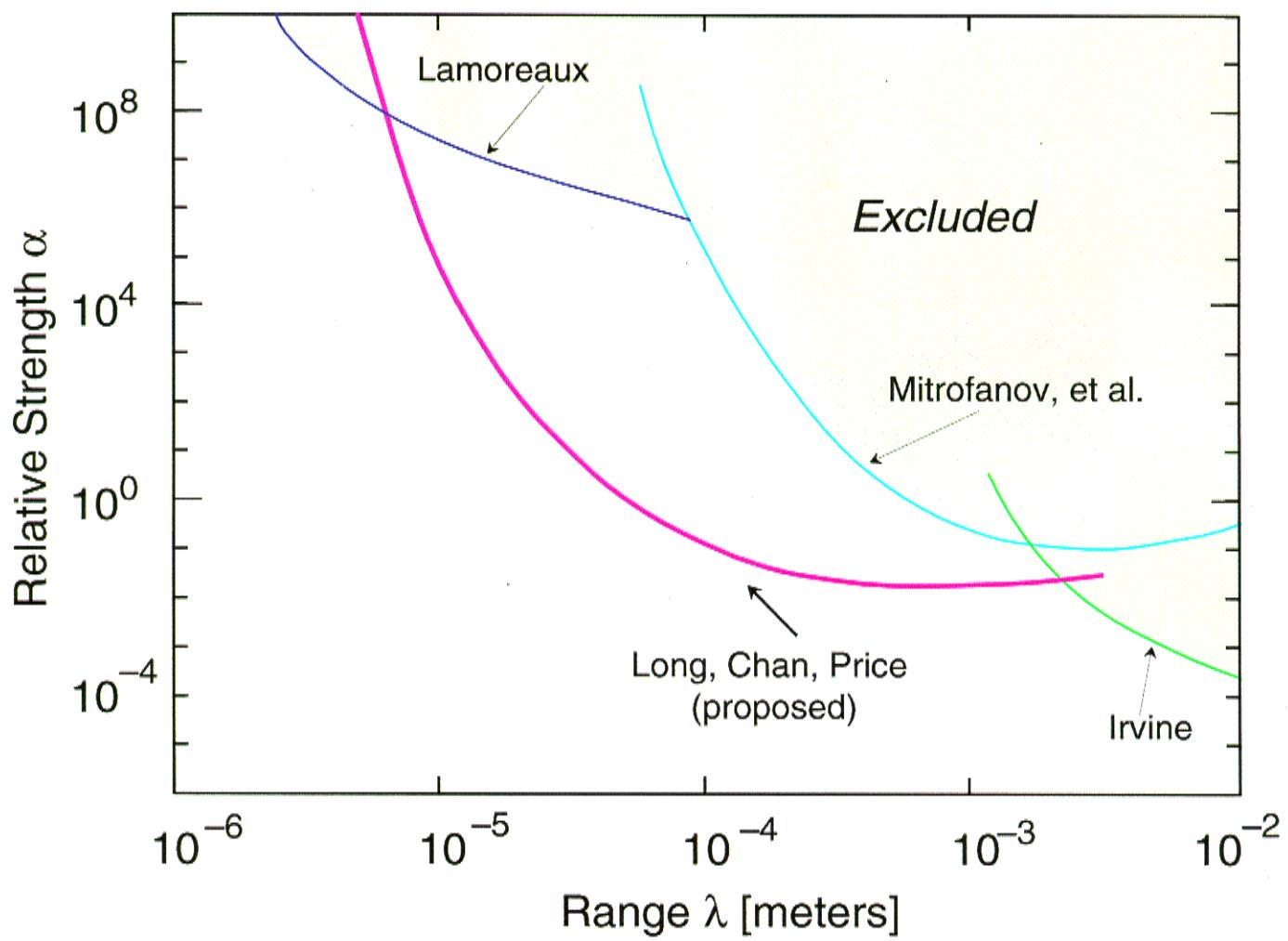
Below 1 mm, limits deteriorate rapidly

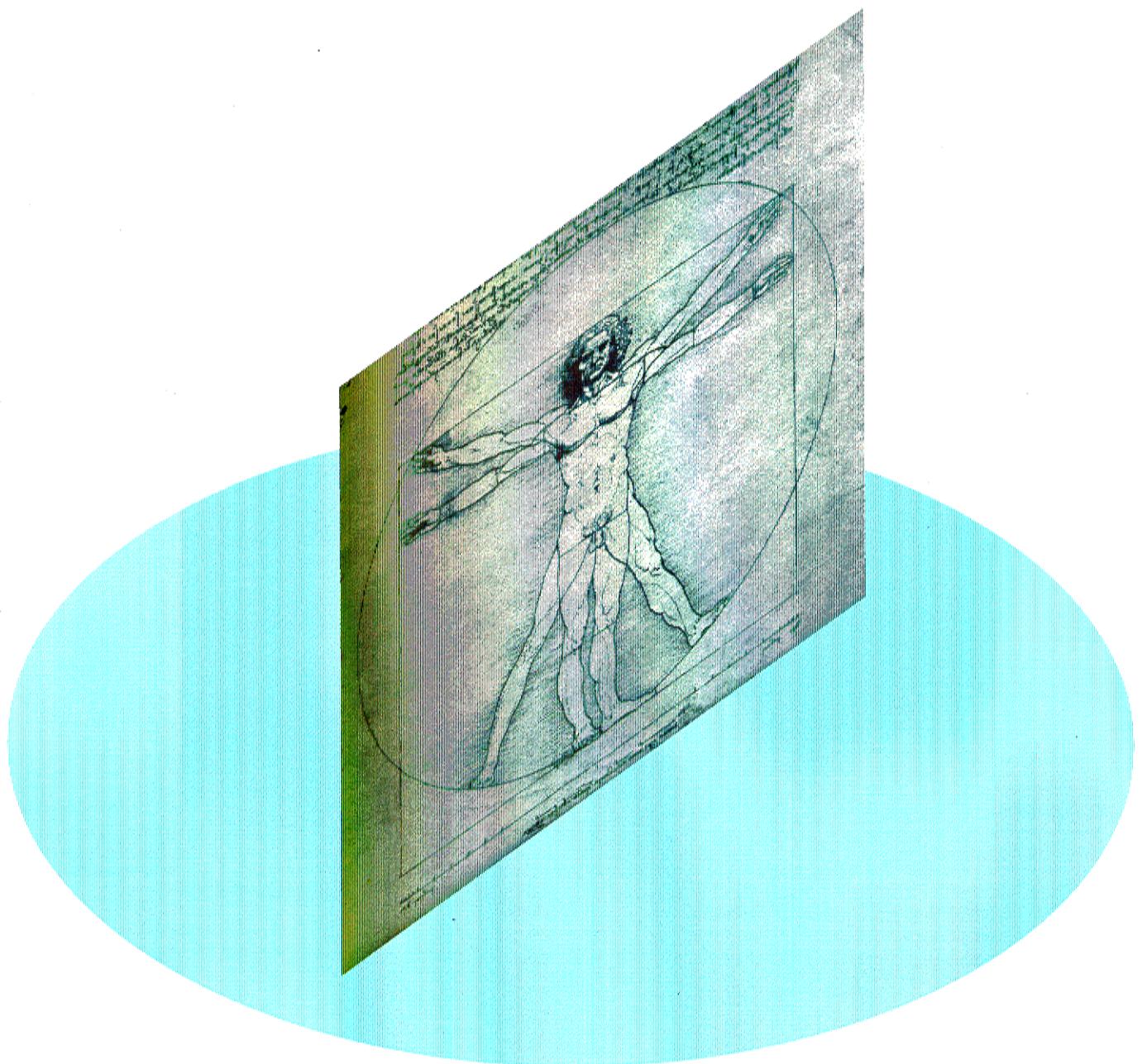
String theory requires 6/7 extra dimensions
Assumed natural to take

$$R_{\text{unobserved}} \approx 1/M_{Pl} \approx 1.6 \times 10^{-35} \text{ m}$$

New Whinkle:

- $SU(3)_c \otimes SU(2)_L \otimes U(1)_Y$ (+ needed extensions)
live on 4-dim branes, not in extra dim.
- Gravity propagates in extra dim.





Dimensional analysis changes: (Gauss's law)

n extra dim. with radius R \Rightarrow

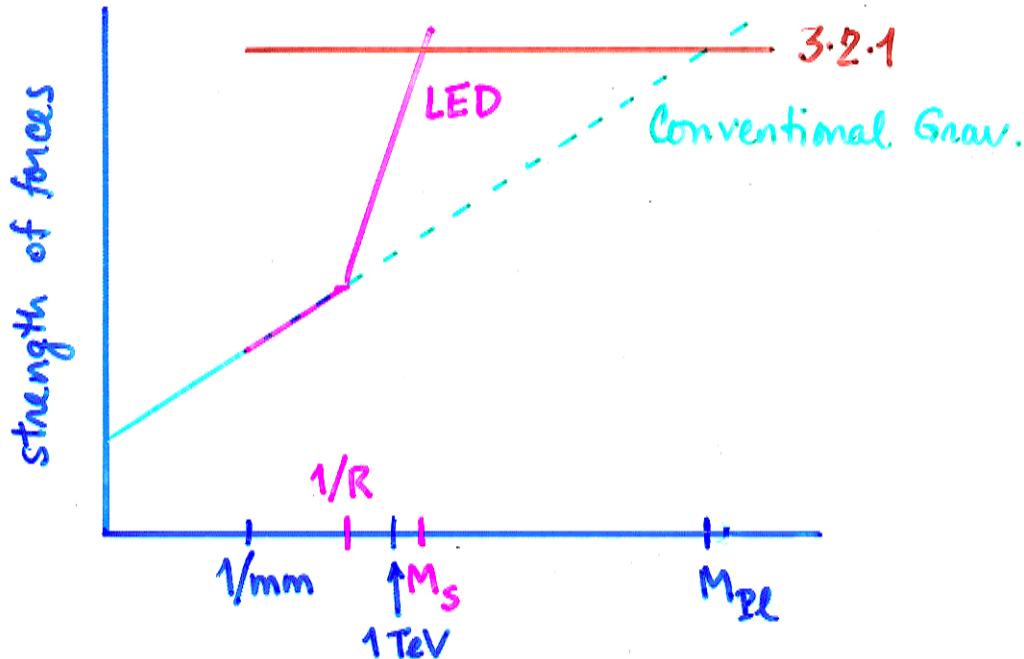
$$G_N \sim M_{Pl}^{-2} \sim M_s^{-n-2} R^{-n}$$

If $M_s \sim 1 \text{ TeV}$, $R \lesssim 1 \text{ mm}$ for $n \geq 2$

$$M_{Pl} = M_s (M_s R)^{n/2}$$

would mean that M_{Pl} results from

a false extrapolation:

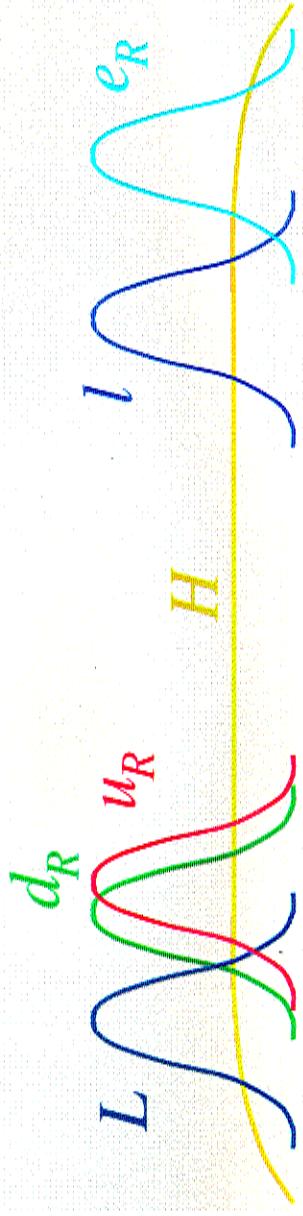


<Gravity's true scale could be near EW scale>

Graviton excitation of a tower of Kaluga-Klein modes in the extra dimensions?

- Missing energy signature for colliders (provatons)

Might Extra Dimensions Explain the Range of Fermion Masses?



$$\mathcal{L}_{\text{Yukawa}} = -\zeta_e [\bar{e}_R (H^\dagger \ell) + (\bar{\ell} H) e_R]; \quad \langle H \rangle_0 = \begin{pmatrix} 0 \\ 1/\sqrt{G_F \sqrt{8}} \end{pmatrix},$$

$$m_e = \zeta_e \cdot 176 \text{ GeV}/c^2 \Rightarrow \zeta_e \approx 3 \times 10^{-6}$$

$$m_t = \zeta_t \cdot 176 \text{ GeV}/c^2 \Rightarrow \zeta_t \approx 1$$

Does localization of fermion wave functions within a brane of small but finite thickness explain the exponential hierarchy?

Micabelli+Schwartz, hep-ph/9912265

Ar Kani-Hamed, Schwartz ...

QCD and the Electroweak Theory

Because QCD is asymptotically free and becomes strong at low energies, it has a rich phase structure...

- Chiral Symmetry Breaking as a (model) mechanism for EWSB.

Ginzburg-Landau \leftrightarrow Higgs (Meissner)

BCS \leftrightarrow XSB (TC)

- Do the other interesting phases of QCD (e.g., color superconductivity) hold lessons for EWSB.

QCD MAY HIDE $SU(2)_L \otimes U(1)_Y$ SYMMETRY

- APPLY $SU(3)_C \otimes SU(2)_L \otimes U(1)_Y$ TO MASSLESS UP + DOWN QUARKS
- TREAT $SU(2)_L \otimes U(1)_Y$ AS PERTURBATION

WITH $m_u = m_d = 0$, \mathcal{L}_{QCD} HAS EXACT (GLOBAL) $SU(2)_L \otimes SU(2)_R$ CHIRAL SYMMETRY AT ENERGY SCALE Λ_{QCD} , STRONG COLOR FORCES BREAK $SU(2)_L \otimes SU(2)_R \rightarrow SU(2)_V$ (ISO SPIN).

SSB \rightarrow GOLDSTONE BOSONS, MASSLESS PIONS.

BROKEN GENERATORS ARE 3 AXIAL CURRENTS COUPLING TO PIONS $\sim f_\pi$.

WHEN $SU(2)_L \otimes U(1)_Y$ IS TURNED ON, EW GAUGE BOSONS COUPLE TO AXIAL CURRENTS AND ACQUIRE MASSES

$$m^2 = \begin{pmatrix} W^+ & W^- & W_3 & a \\ g^2 & 0 & 0 & 0 \\ 0 & g^2 & 0 & 0 \\ 0 & 0 & g^2 & gg' \\ 0 & 0 & gg' & g^2 \end{pmatrix} f_\pi^2 / 4$$

SAME STRUCTURE AS IN STANDARD EW THEORY

DIAGONALIZE:

$$M_W^2 = g^2 f_\pi^2 / 4$$

$$M_Z^2 = (g^2 + g'^2) f_\pi^2 / 4$$

SO THAT

$$M_Z^2 / M_W^2 = (g^2 + g'^2) / g^2 = 1/\cos^2 \theta_W \quad \checkmark$$

$$M_\gamma^2 = 0 \quad \checkmark$$

BUT (BECAUSE $f_\pi = 93 \text{ MeV}$)

$$M_W \approx 30 \text{ MeV}$$

$$M_Z \approx 34 \text{ MeV}$$

1/2650 TIMES TRUE MASSES

(IF ONLY WE DIDN'T KNOW f_π)

TRANSCRIPTION TO TECHNICOLOR

MASSLESS u,d QUARKS \rightarrow NEW FERMIONS
"TECHNIFERMIONS"

QCD \rightarrow NEW INTERACTION
"TECHNICOLOR"

FOR WHICH WE CHOOSE SCALE OF INTERACTION
SO THAT

$$f_{\pi} \rightarrow F_{\pi} = v = (G_F \sqrt{2})^{-1/2}$$

Generates correct M_W, M_Z

but produces no Yukawa couplings, so no fermion masses.

SHOWS POSSIBILITY THAT

Gauge-boson masses
&

Fermion masses

HAVE DIFFERENT ORIGINS.

Wonderful opportunities await particle physics over the next decade:

1-TeV scale (LEP \rightarrow Tevatron \rightarrow LHC)

Problem of Identity (B, K, ν)

Probe dimensionality of spacetime?

If we are inventive enough:

Neutrino factory

e⁺e⁻ Linear collider



A REMARKABLE FLOWERING OF
Experimental Particle Physics
and of

Theory engaged with experiment.